

# Fragmentation of fractal random structures

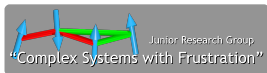
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Institut für Physik, Johannes Gutenberg-Universität Mainz, Germany

with Eren Metin Elçi and Nikolaos G. Fytas (Coventry)

International Conference on Computer Simulation in Physics and Beyond  
Moscow, September 7, 2015

see **Phys. Rev. Lett. 114, 115701 (2015)** and **arXiv:1509.00668**



# Fragmentation



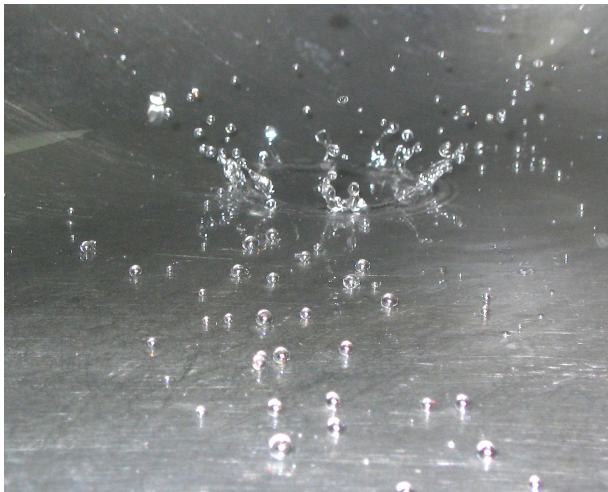
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- $n(s, t)$  number density of  $s$  clusters at time  $t$   
 $a(s)$  fragmentation rate of clusters of size  $s$   
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This leads to a mean-field description with scaling properties.

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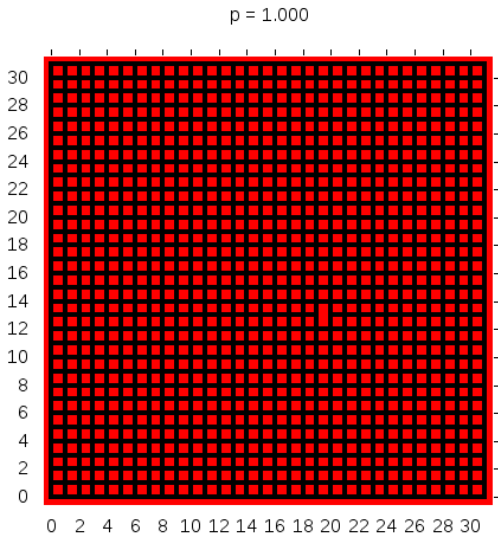
$$Z_{\text{RC}} = \sum_{\mathcal{G}' \subseteq \mathcal{G}} p^{b(\mathcal{G}')} (1-p)^{\varepsilon - b(\mathcal{G}')} q^{k(\mathcal{G}')}, \quad p, q > 0, \quad (2)$$

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- equivalent to Potts model for integer  $q$
- includes percolation ( $q \rightarrow 1$ ), Ising model ( $q = 2$ ), random resistor networks ( $q \rightarrow 0$ ), ...
- continuous percolation transition for  $q < q_c$
- first-order percolation transition for  $q > q_c$
- $q_c = 4$  in 2D,  $q_c \approx 2.8$  in 3D

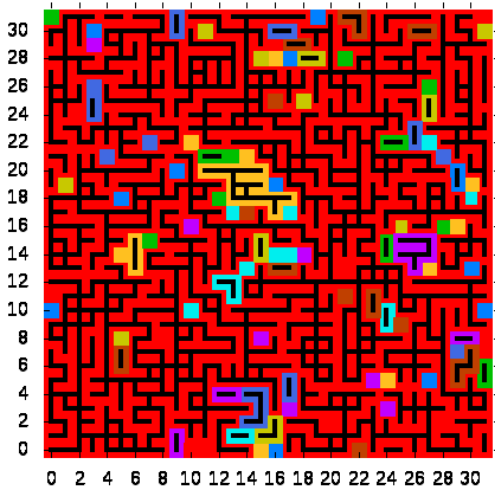
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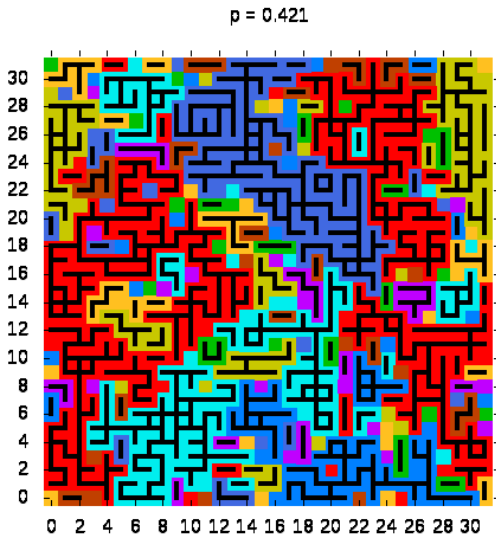


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$p = 0.500$

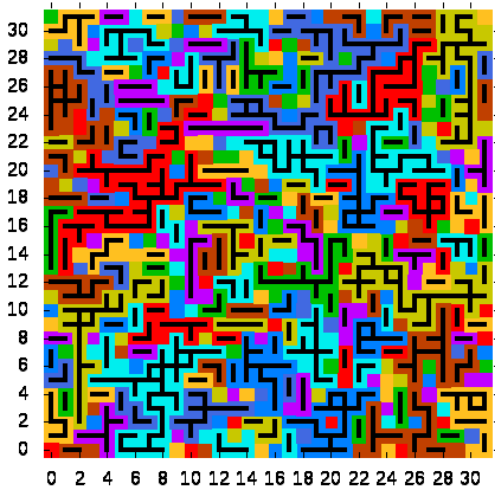


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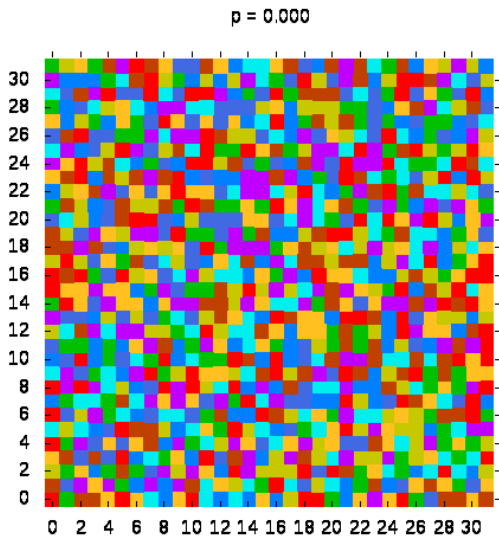


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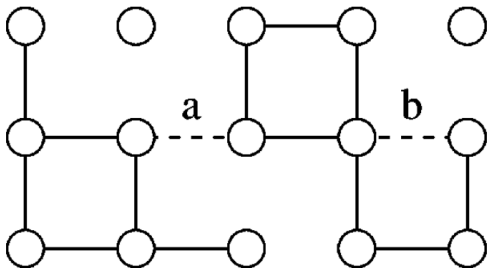


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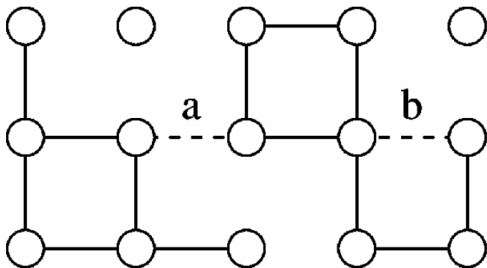
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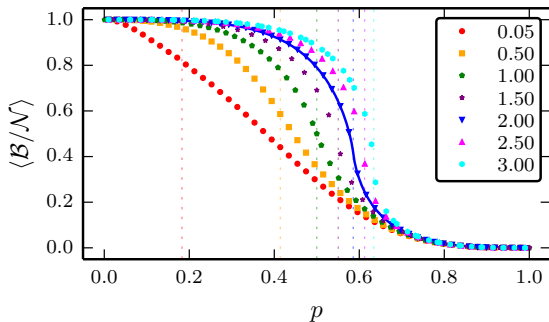
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Fragmentation rate effected by breaking bonds randomly at a constant rate depends on the **density** of bridges.

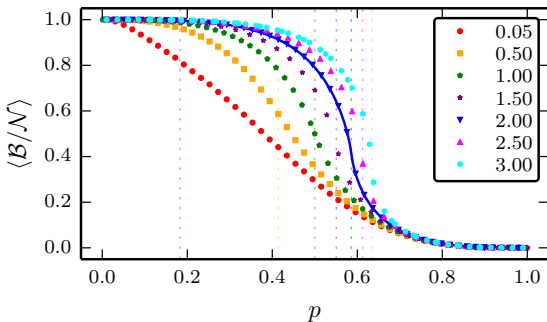
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Critical point  $p_c = \sqrt{q}/(1 + \sqrt{q})$  is where the change in bridge density and hence the **fragility** becomes singular (for  $q \geq 2$ ).



# Scaling theory

**How do the fragments look like?** Scaling theory for  $a(s)$  and  $b(s, s')$  for percolation (Gyure and Edwards, 1992):

$$\begin{aligned} a(s) &\sim s^\lambda \\ b(s, s') &\sim s^{-\phi} g(s'/s) \end{aligned} \tag{3}$$

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Are  $\lambda$  and  $\phi$  related to the known percolation exponents?

## Bridge-edge identity

Based on the Russo-Margulis formula in percolation theory, it is possible to prove the following **bridge-edge identity**

$$\langle \mathcal{B} \rangle = \frac{\langle \mathcal{N} \rangle - p}{(1-p)(1-q)} \quad (5)$$

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Clearly, the case  $q \rightarrow 1$  is singular. There one finds

$$\langle \mathcal{B} \rangle = \frac{1}{p-1} \lim_{q \rightarrow 1} \frac{d}{dq} \langle \mathcal{N} \rangle$$

## Bridge densities

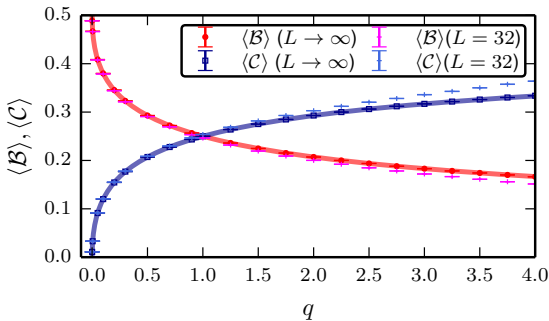
Equation (5) holds for arbitrary graphs. At criticality,  $\langle \mathcal{N} \rangle = 1/2$  and hence

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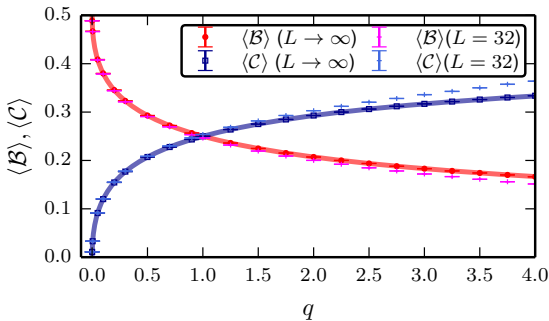
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What are the finite-size corrections?



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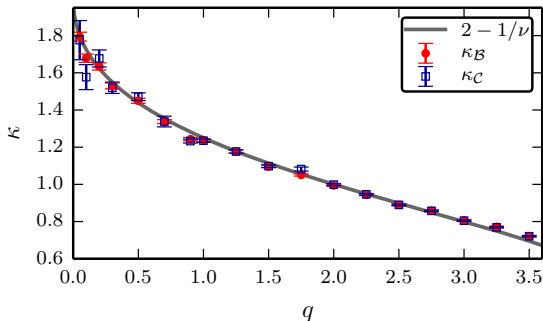
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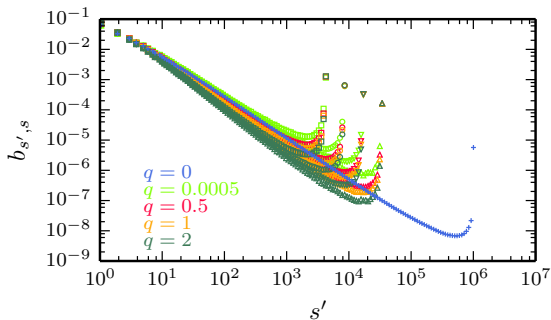


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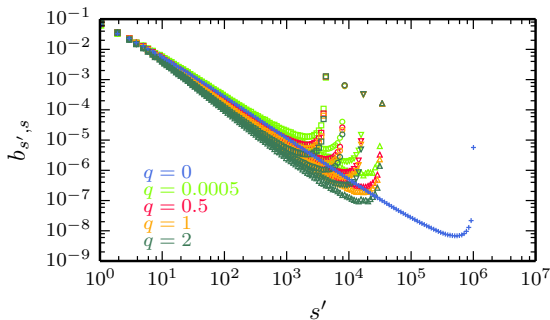
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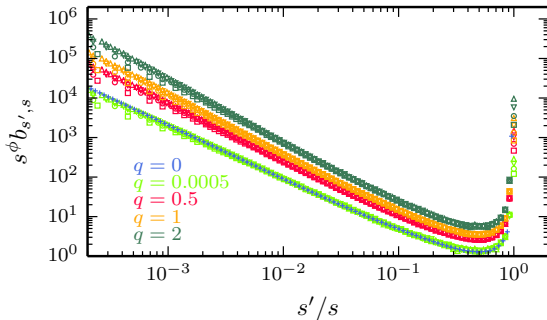


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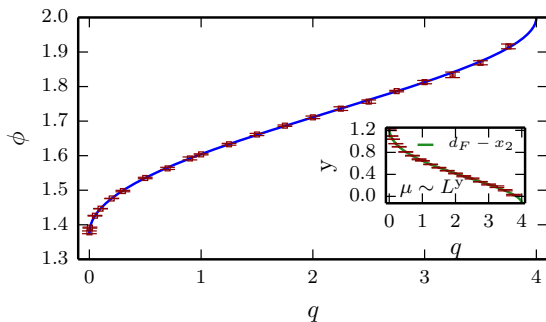
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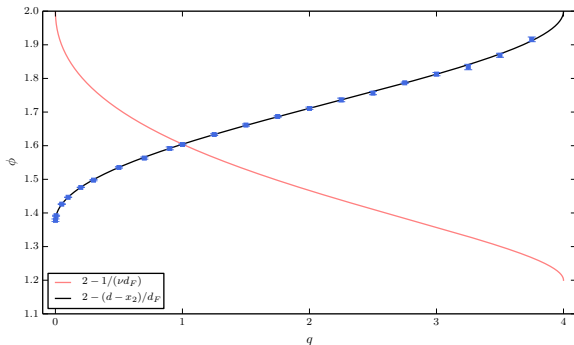
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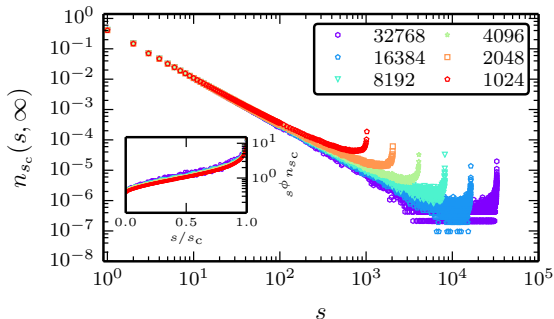
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- introduce **cutoff** to regularize this and model milling or fluid break-up

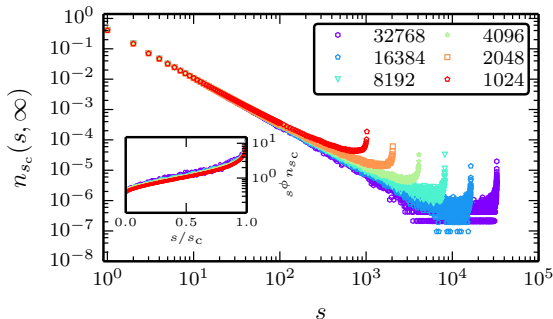
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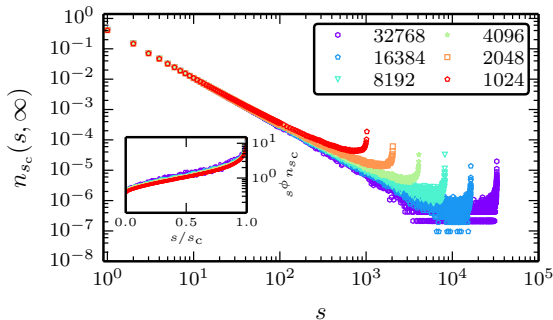
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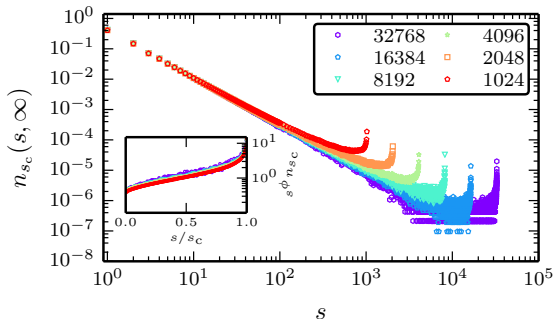


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$$n_{s_c}(s, \infty) \sim s^{-\chi} \mathcal{F}\left(\frac{s}{s_c}\right). \quad (9)$$

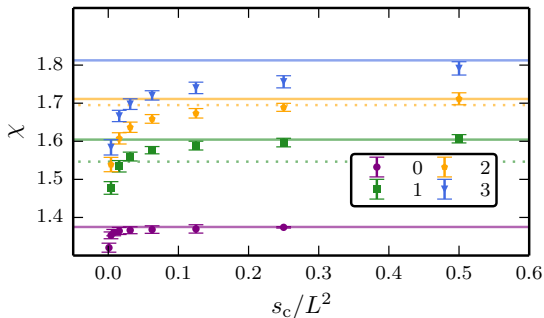
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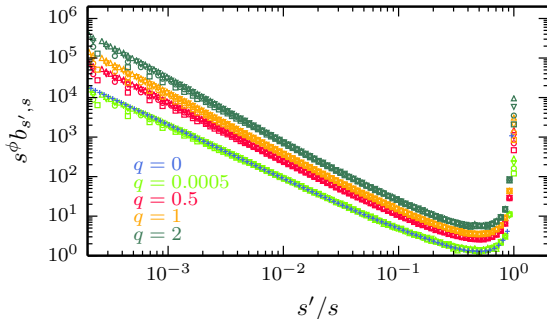
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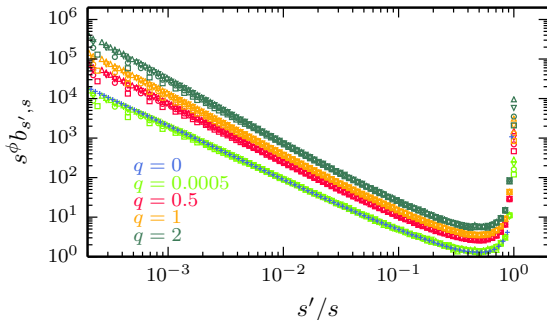
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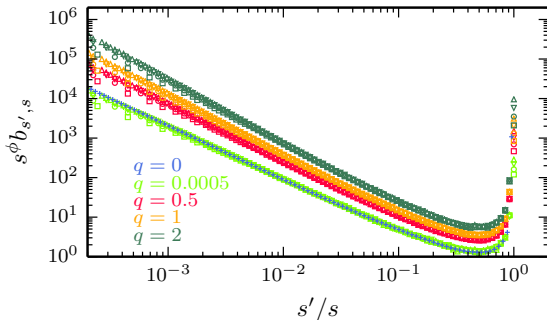


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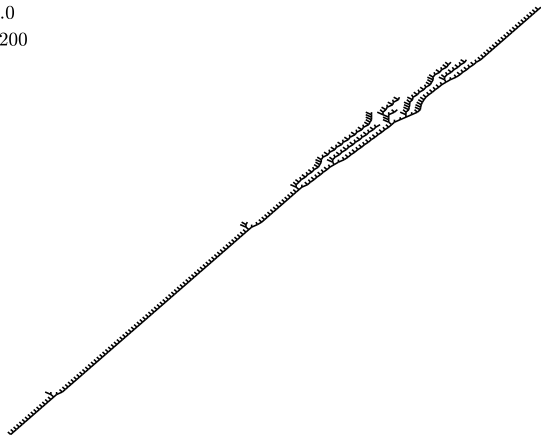
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Genealogical tree of fragmentation events:

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$$q = 1.0$$

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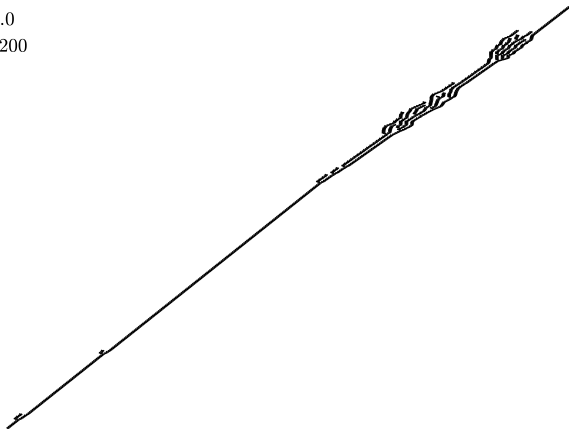
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Genealogical tree of fragmentation events:

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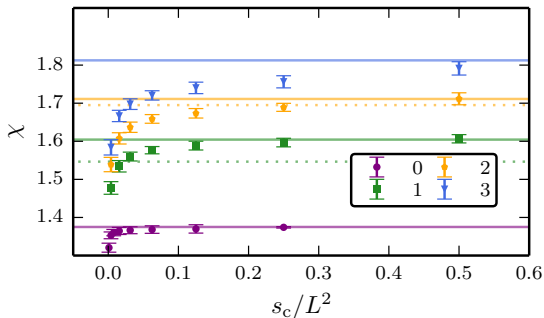
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- can use bridge-free clusters to determine backbone fractal dimension

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# Conclusions

## Summary:

- full characterization of scaling behavior of fragmentation of critical random-cluster model
- new exact relation between bridge and bond densities
- dynamical, non-equilibrium fragmentation with a cutoff leads to power-law fragment distributions, described by equilibrium exponents
- relevant to experiment (universality)?

## Extensions:

- our results extend to three dimensions
- vertex fragmentation related to multi-arm exponents
- can use bridge-free clusters to determine backbone fractal dimension
- dynamical fragmentation of solid objects gives different set of exponents

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